

# Analytical Solutions for the $n$ th Derivatives of Eigenvalues and Eigenvectors for a Nonlinear Eigenvalue Problem

M.S. Jankovic\*

*Esso Resources Canada Limited, Calgary, Canada*

## Abstract

THE nonlinear eigenvalue problem discussed here is defined by the following equations:

$$A(\pi, \lambda) x = 0 \quad (1)$$

$$x^\dagger K(\pi, \lambda) x = 1 \quad (2)$$

where  $A(\pi, \lambda)$  is an  $n \times n$  matrix<sup>†</sup> representing a nonlinear operator on the unrepeat eigenvalue  $\lambda$ ;  $x$  is the corresponding  $1 \times n$  eigenvector that is  $K$ -normalized by an  $n \times n$  positive definite hermitian,  $K$ ;  $x^\dagger$  is the complex conjugate transpose of  $x$ ; and  $\pi$  is a real scalar parameter. The eigenvalue  $\lambda$  is obtained by setting the determinant of  $A$  to zero and solving the resulting characteristic equation for  $\lambda(\pi)$ . The corresponding eigenvector,  $x$ , is then calculated from Eqs. (1) and (2). It is assumed that  $\lambda$  is a differentiable function of  $\pi$ . Elements of  $A$  and  $K$ ,  $a_{ij}$  and  $k_{ij}$  ( $i = 1, \dots, n$ ;  $j = 1, \dots, n$ ) must also be differentiable functions of  $\pi$  and are, in general, nonlinear functions of  $\lambda$  and  $\pi$ .

The objective of this paper is to find explicit analytical solutions for the  $n$ th derivative of the unrepeat eigenvalue  $\lambda$ ,  $\lambda^{(n)}$ , and the  $n$ th derivative of the corresponding eigenvector  $x$ ,  $x^{(n)}$ , with respect to  $\pi$ . Various numerical techniques have been proposed in the past<sup>1</sup> but no analytical solutions of this problem have been available.

## Contents

The differentiation of Eq. (1) with respect to  $\pi$ , with the subscripts  $\pi$  and  $\Pi$  indicating a partial and a total derivative, respectively, gives

$$A_{\Pi} x + A x_{\Pi} = 0 \quad (3)$$

or

$$A_{\lambda} x \lambda_{\Pi} + A x_{\Pi} = -A_{\Pi} x \quad (4)$$

since

$$A_{\Pi} = A_{\pi} + A_{\lambda} \lambda_{\Pi} \quad (5)$$

Similarly, the differentiation of Eq. (2), with respect to  $\pi$  and using

$$K_{\Pi} = K_{\pi} + \lambda_{\Pi} K_{\lambda} \quad (6)$$

gives

$$x^\dagger K_{\lambda} x \lambda_{\Pi} + 2x^\dagger K x_{\Pi} = -x^\dagger K_{\pi} x \quad (7)$$

Equations (4) and (7) form the following system of linear algebraic equations, with  $\lambda_{\Pi}$  and  $x_{\Pi}$  unknown:

$$\left[ \begin{array}{c|c} A & A_{\lambda} x \\ \hline 2x^\dagger K & \alpha \end{array} \right] \begin{pmatrix} x_{\Pi} \\ \lambda_{\Pi} \end{pmatrix} = \begin{pmatrix} -A_{\pi} x \\ -\beta \end{pmatrix} \quad (8)$$

where

$$\alpha = x^\dagger K_{\lambda} x \quad (9)$$

$$\beta = x^\dagger K_{\pi} x \quad (10)$$

The systems of Eq. (8) were obtained, and solved numerically in Ref. 2. The solution of Eq. (8) is written as

$$\begin{pmatrix} x_{\Pi} \\ \lambda_{\Pi} \end{pmatrix} = \left[ \begin{array}{c|c} A & A_{\lambda} x \\ \hline 2x^\dagger K & \alpha \end{array} \right]^{-1} \begin{pmatrix} -A_{\pi} x \\ -\beta \end{pmatrix} \quad (11)$$

Two distinct cases,  $\alpha = 0$  and  $\alpha \neq 0$ , are considered. The solutions for  $\alpha = 0$  are

$$2x^\dagger K x_{\Pi} = -\beta \quad (12)$$

$$x_{\Pi} = -\frac{\beta}{2} x \quad (13)$$

The solution for  $\alpha \neq 0$  is given as

$$\begin{pmatrix} x_{\Pi} \\ \lambda_{\Pi} \end{pmatrix} = \left[ \begin{array}{c|c} P & x/2 \\ \hline -2x^\dagger K P/\alpha & 0 \end{array} \right] \begin{pmatrix} -A_{\pi} x \\ -\beta \end{pmatrix} \quad (14)$$

or

$$x_{\Pi} = -P A_{\pi} x - \frac{\beta}{2} x \quad (15)$$

and

$$\lambda_{\Pi} = 2 x^\dagger \left( K \frac{P A_{\pi}}{\alpha} \right) x \quad (16)$$

Theorems for both  $\alpha = 0$  and  $\alpha \neq 0$  are stated here without proofs, which are given in the unabridged version of this paper.

**Theorem I:** For the nonlinear eigenvalue problem previously defined, and for  $\alpha = 0$

$$-(A_{\lambda}^{-1} A_{\pi})^{(n)} x = \lambda^{(n+1)} x$$

The superscripts  $n$  and  $n+1$  indicate the  $n$ th and  $n+1$ st derivatives with respect to the parameter  $\pi$ .

**Theorem II:** For the nonlinear eigenvalue problem previously defined, and for  $\alpha = 0$ , the  $n$ th derivative of the eigenvalue  $\lambda$  with respect to the parameter  $\pi$  is

$$\lambda^{(n)} = -x^\dagger K (A_{\lambda}^{-1} A_{\pi})^{(n-1)} x$$

**Theorem III:** For the nonlinear eigenvalue problems previously defined, and for  $\alpha = 0$ , the  $n$ th derivative of the eigenvector  $x$  is

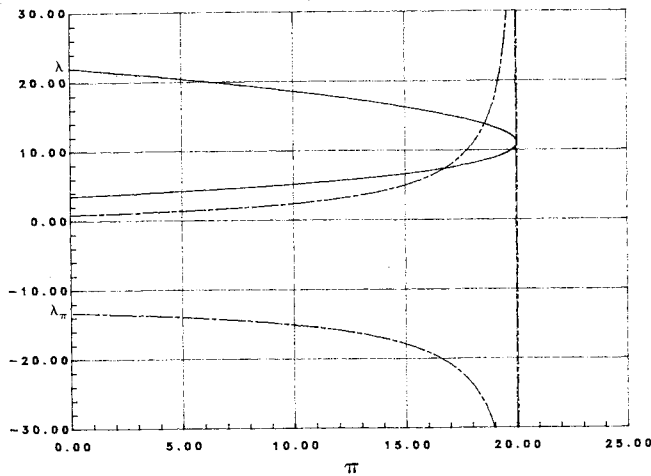
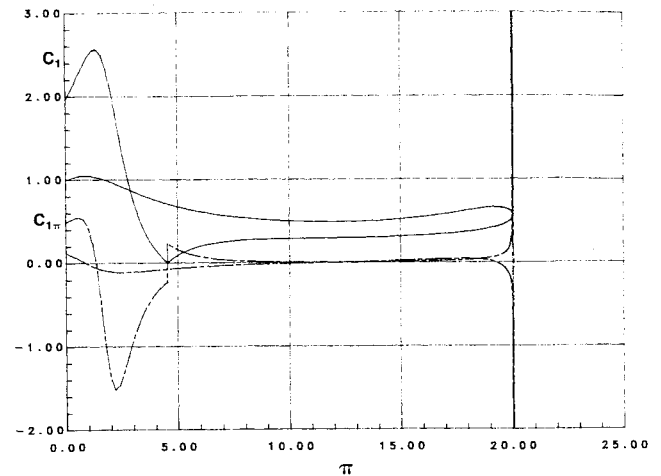
$$x^{(n)} = -\frac{1}{2} (\beta x)^{(n-1)}$$

**Theorem IV:** For the nonlinear eigenvalue problem previously defined, and for  $\alpha = 0$

$$A^{(n-k)} x^{(k)} = 0, \quad k = 0, 1, \dots, n$$

Received Aug. 24, 1987. Copyright © American Institute of Aeronautics and Astronautics, Inc., 1987. All rights reserved. Full paper available from National Technical Information Service, Springfield, VA 22151, at the standard price (available upon request).

\*Research Specialist, Research Department.

Fig. 1  $\lambda$  and  $\lambda_\pi$ .Fig. 2  $C_1$  and  $C_{1\pi}$ .

**Theorem V:** For the nonlinear eigenvalue problem previously defined, and for  $\alpha \neq 0$ , the  $n$ th derivative of the eigenvalue  $\lambda$  with respect to the parameter  $\pi$  is

$$\lambda^{(n)} = 2 \left[ x^\dagger K \left( \frac{P A_\pi}{\alpha} \right) x \right]^{(n-1)}$$

**Theorem VI:** For the nonlinear eigenvalue problem previously defined, the  $n$ th derivative of the eigenvector  $x^{(n)}$  with respect to the parameter  $\pi$  is

$$x^{(n)} = -(PA x)_\pi^{(n-1)} - \frac{(\beta x)^{(n-1)}}{2}$$

These theorems are also valid for the linear eigenvalue problem.

The theory outlined here applies to many problems in dynamics, and other branches of applied mathematics such as the theories of oscillation, elasticity, and optimization. Its application to an elementary, nonconservative problem of the theory of elasticity, namely the problem of the oscillating bar compressed by a tangential force, follows. In Ref. 3, the matrix  $A$  is given by

$$\begin{aligned} a_{11} &= a_{13} = 0 \\ a_{22} &= a_{24} = 0 \\ a_{33} &= a_{34} = 0 \\ a_{43} &= a_{44} = 0 \\ a_{12} &= a_{14} = 1 \\ a_{21} &= \beta \\ a_{23} &= \hat{\alpha} \\ a_{31} &= \beta^2 \sin \hat{\beta} + \hat{\alpha} \beta \sinh \hat{\alpha} \\ a_{32} &= \beta^2 \cos \hat{\beta} + \hat{\alpha}^2 \cosh \hat{\alpha} \\ a_{41} &= -(\beta^3 \cos \hat{\beta} + \beta \hat{\alpha}^2 \cosh \hat{\alpha}) \\ a_{42} &= \beta^3 \sin \hat{\beta} - \hat{\alpha}^3 \sinh \hat{\alpha} \end{aligned} \quad (17)$$

which was obtained by substituting the solution  $V(\xi) = c_1 \sin \beta \xi + c_2 \cos \beta \xi + c_3 \sinh \hat{\alpha} \xi + c_4 \cosh \hat{\alpha} \xi$  into the governing equation

$$\frac{d^4 V}{d\xi^4} + \pi \frac{d^2 V}{d\xi^2} - \omega^2 V = 0, \quad \xi = \frac{Z}{\ell} \quad (18)$$

and by using the appropriate boundary conditions. In Eq. (18),  $\omega$  is a nondimensional vibration frequency, and the nondimensional parameter  $\pi$  is proportional to the tangential force  $P$ , is the square of the bar's length  $\ell$ , and is inversely proportional to the bending stiffness  $EJ$ .

That is

$$\pi = \frac{P \ell^2}{EJ} \quad (19)$$

Also,  $\hat{\alpha}$  and  $\hat{\beta}$  are defined as

$$\hat{\alpha} = \sqrt{(\pi/4 + \lambda) 1/2 - \pi/2^{1/2}} \quad (20a)$$

$$\hat{\beta} = \sqrt{(\pi/4 + \lambda) 1/2 - \pi/2^{1/2}} \quad (20b)$$

where  $\lambda \equiv \omega^2$  and  $\omega$  was calculated from  $\det A = 0$  as a function of  $\pi$ . The normality condition

$$\int_0^1 V^2 d\xi = 1 \quad (21)$$

can be written as

$$x^\dagger K x = 1 \quad (22)$$

where the elements of the matrix  $K$  are evaluated in closed form and the eigenvector is given by  $x^\dagger = (c_1 \ c_2 \ c_3 \ c_4)$ . The critical value of the parameter  $\pi$ ,  $\pi^*$ , which gives the critical frequency, is calculated from Eq. (16) to be  $\pi^* = 20.05059 \dots$ . The derivatives of  $\lambda$  and  $c_1$  are shown in Figs. 1 and 2.

The theory presented here was used to calculate the first and second derivatives of the solar array's eigenvalues and eigenvectors for the Hermes satellite. The problem is quite complex and is discussed in Ref. 4, and in detail in Ref. 2.

## Conclusions

This paper analytically calculates the  $n$ th derivatives of eigenvalues and eigenvectors for a nonlinear eigenvalue problem. It also states the related theorems. Some applications are mentioned and an example is discussed.

## Acknowledgment

I wish to thank Professor Peter Lancaster from the University of Calgary and Dr. A.L. Andrew from La Trobe University for their helpful comments.

## References

- Andrew, A.L., "Eigenvalue Problems with Nonlinear Dependence on the Eigenvalue Parameter," Bibliography, Tech. Rept., Mathematics Dept., La Trobe Univ., 1974.
- Jankovic, M.S., "Deployment Dynamics of Flexible Spacecrafts," Appendix A, Ph.D. Thesis, Univ. of Toronto, 1979, p. 42-45.
- Bolotin, V.V., *Nonconservative Problems of the Theory of Elastic Stability*, Pergamon Press, 1963, p. 90-93.
- Jankovic, M.S., Comments on "Dynamics of a Spacecraft during Extension of Flexible Appendages," *Journal of Guidance, Control, and Dynamics*, Vol. 7, Jan.-Feb. 1984, p. 128.